THE ASTROPHYSICAL JOURNAL, 181:L75-L77, 1973 April 15 © 1973. The American Astronomical Society. All rights reserved. Printed in U.S.A.

THE REFLECTION EFFECT IN HZ HERCULIS

ROBERT E. WILSON

Institute for Space Studies, New York, N. Y. Received 1972 December 11; revised 1973 February 2

ABSTRACT

The theoretical phase law for the reflection effect is not in agreement with the photometric observations of HZ Her. The discrepancy may be explained by transfer of energy from the irradiated side of the optical component to the side facing away from the X-ray source. This transfer must involve a significant fraction of the envelope of the optical component, and is in accord with recently developed ideas on the behavior of irradiated convective envelopes.

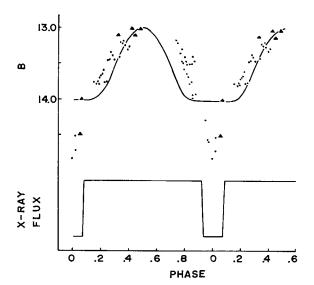
Subject headings: binaries — X-ray sources

Recently Davidsen et al. (1972) and Forman, Jones, and Liller (1972) have published photometric observations of the X-ray eclipsing system HZ Herculis (= Her X-1). Forman et al. interpreted the 1.5 mag brightness variation as due to the usual binary-star reflection effect taking place on the relatively large optical component, with the source of energy being the compact X-ray source. While it seems scarcely conceivable that the variation is due to something other than reflection, there does exist a major difficulty for straightforward acceptance of this idea. The problem is that the theoretical phase law for the "reflected" light from a rotating binary component is so completely at variance with the observed light curves that the discord can easily be seen even from the photographic observations of Forman et al. A convenient summary of early theoretical work on the phase law for reflection can be found in Russell and Merrill (1952). A law by Milne (1926), as modified by Sen (1948) to account for the finite distance of the energy source, is adequate for HZ Her, where the irradiating energy is emitted by a point source. Assuming a value of 0.4 for r_1/a , the radius of the optical component relative to the component separation, the Milne-Sen phase law for orbital inclination 90° is

$$f(\theta) = 0.314 + 0.447 \cos \theta + 0.153 \cos 2\theta + 0.005 \cos 3\theta - 0.006 \cos 4\theta$$
.

Checks made with the binary-star program of Wilson and Devinney (1971) show that this relation is not modified very much by tidal distortion or by inclinations less than 90° , provided that i is not less than about 65° . Of course, the constant term will be larger than 0.314 if there are constant sources of light in the system, and a value of 0.9 was used in the present computations.

Figure 1 compares the Milne-Sen law with the (mostly) photoelectric B observations by Davidsen et al. The flat region in the Milne-Sen law is due to the fact that only the irradiated region of the star has a strong brightness variation, and this part is out of sight when we view the dark side. The observations show an effect quite opposite to that of the theoretical law—there is a sharp downward spike at just the phase where there should be a broad minimum, while the maximum is much more gently curved than the theoretical law. The photographic observations (Forman et al.) also differ clearly from the theoretical law, but because of their scatter the disagreement is somewhat less striking. Remarkably, the Milne-Sen law agrees almost reasonably well with the observations if it is shifted 180° and inverted (turned upside down)! Since the observed light curves show all the gross features of the reflection effect (maxima at X-ray phase 0.5, minima at phase 0.0; approximate symmetry about phases 0.0 and 0.5) and the spectral type and color index vary as expected, it seems certain that the basic interpre-



Vol. 181

Fig. 1.—The expected phase law for a normal reflection effect (solid curve) compared with the B light curve of HZ Her by Davidsen et al. Most observations are photoelectric (dots). A few photographic observations are shown as triangles.

tation by Forman et al. is correct. Nevertheless, it is entirely impossible for the reflection effect, as encountered in normal binaries, to show the phase dependence observed for HZ Her.

We might try to explain the spike at minimum as due to the eclipse of a large, optically emitting region surrounding the X-ray star. However, this idea cannot explain the failure of the Milne-Sen law to fit the observations at other phases, such as the maximum. At first sight it would seem that the improved agreement found upon inverting the theoretical phase law is merely coincidental and should not be given any particular interpretation. However, it is instructive to ask what geometrical situation corresponds to turning the phase law upside down. This operation provides the phase law for the photographic negative of a motion picture of the rotating component. That is, if we replace each (relatively) bright element of area on a normal reflecting star with an element which is relatively dark, and vice versa, the inverted Milne-Sen law would be a good approximation to the resulting light curves. Now an ordinary binary component which shows a strong reflection effect has a very bright cap whose intensity diminishes rapidly for points successively farther from the substellar point. For practical purposes it can be said that the bright region extends substantially less than halfway around the star, since both the inverse square law and a projection factor greatly diminish the heating effect at the "sides" of the star. The primary of HZ Her behaves as if it were the negative image of such a normal reflecting component. Therefore, we expect it to have a very dark cap on one end which extends less than halfway around the star. If this dark cap in interpreted as the anti-substellar point, it appears that the actual reflection effect extends substantially more than halfway around the star, a situation which would seem to be geometrically impossible since the back hemisphere is not in view of the X-ray source. We are forced to conclude that the large X-ray flux drives circulation currents which redistribute the incident energy over a considerable part of the surface which is not exposed to direct X-radiation. Furthermore, these circulation currents cannot be superficial surface currents, since the characteristic time for radiation of the heat content of the surface layers is not longer than a few minutes, even under the most extreme assumptions, and this time scale would not allow adequate time for transfer of the energy over the surface. We conclude that a major fraction of the envelope must be affected by the X-ray flux, that circulation currents carry energy into the interior, and that this energy subsequently appears at points of the surface which are more than halfway around the star from the substellar point. Forman et al. have estimated the undisturbed spectral type of the optical component to be middle F, and this estimate would be too early if some of the back hemisphere is effectively heated by the X-rays. Therefore, we expect a deep convective envelope for this star. It has already been shown by Rucinski (1969; see also Lucy 1968) that for stars with fully adiabatic convective envelopes, the bolometric albedo should be only about 0.4–0.5 rather than unity, because the incident flux disturbs the envelope to quite deep layers. Such values for the bolometric albedo have been found from observations of several binaries (e.g., Napier 1970; Wilson et al. 1972). The energy which is not "reflected" locally must appear at other points of the surface. HZ Herculis may be an unusually favorable case for testing such ideas, and it should be most useful to construct a model for its interior structure.

Note added 1972 December—New observations by Bahcall and Bahcall (1972) have just appeared. The Bahcalls fitted their observations on the basis of a simple model of a spherical component which is uniformly bright over one hemisphere and uniformly dim over the other. Their model neglects two important effects, which are (1) that the irradiated cap is less than a hemisphere due to the finite (in fact, small) distance of the X-ray source; (2) that the X-rays are diluted by the surface (slant) projection factor and by the inverse square law. Since we have verified by detailed computations (including cases for $i \neq 90^{\circ}$) that a satisfactory match between the observations and the light curve for a normal reflection effect is not possible if one accounts for (1) and (2) above, it seems clear that the agreement shown by figure 2 of Bahcall and Bahcall should be regarded as fortuitous.

REFERENCES

Bahcall, J. N., and Bahcall, N. A. 1972, Ap. J. (Letters), 178, L1.
Davidsen, A., Henry, J. P., Middleditch, J., and Smith, H. E. 1972, Ap. J. (Letters), 177, L97.
Forman, W., Jones, C. A., and Liller, W. 1972, Ap. J. (Letters), 177, L103.
Lucy, L. B. 1968, Ap. J., 153, 877.
Milne, E. A. 1926, M.N.R.A.S., 87, 43.
Napier, W. M. 1970, Ap. and Space Sci., 11, 475.
Rucinski, S. M. 1969, Acta Astr., 19, 245.
Russell, H. N., and Merrill, J. E. 1952, Contr. Princeton Univ. Obs., 26, 44–45.
Sen, H. K. 1948, Proc. Nat. Acad. Sci., 34, 311.
Wilson, R. E., Deluccia, M. R., Johnston, K., and Mango, S. A. 1972, Ap. J., 177, 191.
Wilson, R. E., and Devinney, E. J. 1971, Ap. J., 166, 605.